

4.7 We are given a ton of classical diatomic molecules in a volume

$$V \text{ and temperature } T \text{ w/ Hamiltonian: } H(F_1, \bar{F}_2, \bar{P}_1, \bar{P}_2) = \frac{1}{2m}(P_1^2 + P_2^2) + \frac{1}{2}k_0|\vec{F}_1 - \vec{F}_2|^2$$

[NOTE: I changed capital K to k₀ for the spring constant for clarity.]

We'll begin by writing the 1-particle partition function Q₁.

$$Q_1 = \int d\omega e^{-\beta E}$$

In this case, we're using a 6-D phase space, so $d\omega = \frac{d^6 r d^6 p}{h^6}$

$$Q_1 = \iint d^6 r d^6 p e^{-\beta \frac{P_1^2}{2m}} e^{-\beta \frac{P_2^2}{2m}} e^{-\beta \frac{1}{2}k_0|\vec{F}_1 - \vec{F}_2|^2}$$

Noting that the volume is pretty big, we can pretend it's infinite, and the momentum integrals become easy (just Gaussian integrals).

$$(d^6 p = dP_{1x} dP_{1y} dP_{1z} dP_{2x} dP_{2y} dP_{2z}, \text{ and } P_i^2 = P_{ix}^2 + P_{iy}^2 + P_{iz}^2, \text{ etc.})$$

$$Q_1 = \frac{1}{h^6} \left(\frac{2\pi m}{\beta} \right)^{6/2} \int d^6 \vec{r} e^{-\frac{\beta}{2}k_0|\vec{F}_1 - \vec{F}_2|^2}$$

Next, let $\vec{x} = \vec{F}_1 - \vec{F}_2$, and $\vec{y} = \vec{F}_2 \Rightarrow \vec{F}_1 = \vec{x} + \vec{y}$ and $\vec{F}_2 = \vec{y}$

This allows us to change variables in our remaining integral.

$$\rightarrow |\vec{F}_1 - \vec{F}_2|^2 = \vec{x}^2$$

$$\rightarrow d^6 \vec{r} = d^3 \vec{F}_1 d^3 \vec{F}_2 = d^3 \vec{x} d^3 \vec{y} \underbrace{\frac{\partial(\vec{F}_1, \vec{F}_2)}{\partial(\vec{x}, \vec{y})}}_{= d^3 \vec{x} d^3 \vec{y}}$$

This is the Jacobian Determinant. I calculated in Mathematica that it simply equals 1. Nice!

$$Q_1 = \frac{1}{h^6} \left(\frac{2\pi m}{\beta} \right)^{6/2} \int d^3 \vec{x} d^3 \vec{y} e^{-\frac{\beta}{2}k_0 \vec{x}^2} = \frac{V}{h^6} \left(\frac{2\pi m}{\beta} \right)^{6/2} \int d^3 \vec{x} e^{-\frac{\beta}{2}k_0 \vec{x}^2}$$

This is 3 identical integrals multiplied. we get

$$Q_1 = \frac{V}{h^6} \left(\frac{2\pi m}{\beta} \right)^{6/2} \left(\frac{2\pi}{k_0 \beta} \right)^{3/2} = V \left(\frac{2\pi m}{h^2 \beta} \right)^3 \left(\frac{2\pi}{k_0 \beta} \right)^{3/2} = V f(T)$$

(where $f(T) = \frac{(2\pi k T)^{9/2}}{h^6} \left(\frac{m^2}{k_0} \right)^{3/2}$)

Finally, $\boxed{Q_1(z, V, T) = e^{zVf(T)}}$ by 4.4.3

Thermodynamics (by 4.4.5-8, we use f(T):)

$$P = z k T f(T) \Rightarrow \boxed{P = e^{zVf(T)} k T \left(\frac{(2\pi k T)^{9/2}}{h^6} \left(\frac{m^2}{k_0} \right)^{3/2} \right)^{1/2}} \quad N = z V f(T) \Rightarrow \boxed{N = V B P}$$

$$U = z V k T^2 f'(T) \Rightarrow U = \frac{9}{2} e^{\frac{4}{3} k T} V k T \left(2\pi k T \right)^{9/2} \left(\frac{m^2}{k_0} \right)^{3/2} \Rightarrow \boxed{U = \frac{9}{2} V P = \frac{9}{2} N k T}$$

$$A = N k T \ln(z) = z V k T f'(T) = N k T \left(\frac{u}{k T} \right) - PV \Rightarrow \boxed{A = Nu - PV}$$

$$S = -N k \ln(z) + z V k T f'(T) + z V k f(T) = -\frac{N k u}{k T} + \frac{U}{T} + \frac{PV}{T}$$

$$\Rightarrow \boxed{S = \frac{1}{T}(U + PV - Nu)}$$

4.7, cont

We have called $\vec{F}_{12} = \vec{x}$.

We are looking for $\langle x^2 \rangle$

Using equation ② on the previous page, we use

$$\begin{aligned}\langle x^2 \rangle &= \frac{\frac{V}{h^3} \left(\frac{2\pi m}{\beta}\right)^{6/2} \int d^3 \vec{x} x^2 e^{-\frac{\beta k_0}{2} x^2}}{\frac{V}{h^3} \left(\frac{2\pi m}{\beta}\right)^{6/2} \int d^3 \vec{x} e^{-\frac{\beta k_0}{2} x^2}} \\ &= \frac{\int dx_1 dx_2 dx_3 (x_1^2 + x_2^2 + x_3^2) e^{-\frac{\beta k_0}{2} (x_1^2 + x_2^2 + x_3^2)}}{\int dx_1 dx_2 dx_3 e^{-\frac{\beta k_0}{2} (x_1^2 + x_2^2 + x_3^2)}} \quad \text{See Mathematica} \\ &= \frac{6}{\beta k_0} = \frac{3kT}{k_0}\end{aligned}$$

$$\Rightarrow \boxed{\langle F_{12}^2 \rangle = \frac{3kT}{k_0}}$$

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This is the jacobian matrix.

```
x = {x1, x2, x3};
y = {y1, y2, y3};
r1 = x + y;
r2 = y;
Det[{{D[r1[[1]], x[[1]]], D[r1[[1]], x[[2]]]}, {D[r1[[1]], x[[3]]], D[r1[[1]], y[[1]]], D[r1[[1]], y[[2]]], D[r1[[1]], y[[3]]]}, {D[r1[[2]], x[[1]]], D[r1[[2]], x[[2]]], D[r1[[2]], x[[3]]], D[r1[[2]], y[[1]]]}, {D[r1[[2]], y[[2]]], D[r1[[2]], y[[3]]]}, {D[r1[[3]], x[[1]]], D[r1[[3]], x[[2]]], D[r1[[3]], x[[3]]], D[r1[[3]], y[[1]]]}, {D[r1[[3]], y[[2]]], D[r1[[3]], y[[3]]]}, {D[r2[[1]], x[[1]]], D[r2[[1]], x[[2]]], D[r2[[1]], x[[3]]], D[r2[[1]], y[[1]]]}, {D[r2[[1]], y[[2]]], D[r2[[1]], y[[3]]]}, {D[r2[[2]], x[[1]]], D[r2[[2]], x[[2]]], D[r2[[2]], x[[3]]], D[r2[[2]], y[[1]]]}, {D[r2[[2]], y[[2]]], D[r2[[2]], y[[3]]]}, {D[r2[[3]], x[[1]]], D[r2[[3]], x[[2]]], D[r2[[3]], x[[3]]], D[r2[[3]], y[[1]]]}, {D[r2[[3]], y[[2]]], D[r2[[3]], y[[3]]]}]
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This is the integral on the final page. I have replaced the x's with a, b, and c.

$$\left(\int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} (a^2 + b^2 + c^2) e^{-\frac{\beta k_0}{2} (a^2 + b^2 + c^2)} da \right) db \right) dc \right) / \left(\int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} e^{-\frac{\beta k_0}{2} (a^2 + b^2 + c^2)} da \right) db \right) dc \right)$$

ConditionalExpression[$\frac{3}{k_0 \beta}, (\text{Re}[k_0] \neq 0 \text{ || } k_0 \notin \text{Reals}) \& \& \text{Re}[k_0 \beta] > 0 \& \& \text{Re}[k_0 \beta] > 0$]