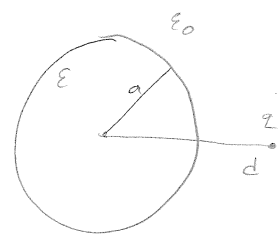


#3

Jackson 4.9(a)



Outside: Use images. Put q' at $y = a^2/d$ as on hub 2,

$$\text{So } \Phi_{\text{out}} = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{|\vec{x}-d\hat{y}|} + \frac{q'}{|\vec{x}-\left(\frac{a^2}{d}\right)\hat{y}|} \right)$$

and use the fact that we have spherically symmetric B.C.'s to

$$\text{get } \frac{1}{|\vec{x}-\vec{x}'|} = \sum_{l=0}^{\infty} \frac{r_l^l}{r'^{l+1}} P_l(\cos\theta)$$

$$\text{Thus } \Phi_{\text{out}}(r, \theta) = \frac{1}{4\pi\epsilon_0} \left(\sum_{l=0}^{\infty} \frac{q d^l}{r^{l+1}} P_l(\cos\theta) + \frac{q' \left(\frac{a^2}{d}\right)^l}{r^{l+1}} P_l(\cos\theta) \right) \quad |\vec{x}| > d$$

$$\rightarrow \Phi_{\text{out}}(r, \theta) = \frac{1}{4\pi\epsilon_0} \left(\sum_{l=0}^{\infty} \frac{P_l(\cos\theta)}{r^{l+1}} \left(q d^l + q' \frac{a^{2l}}{d^l} \right) \right) \quad |\vec{x}| > d \quad \checkmark$$

$$\rightarrow \text{and } \Phi_{\text{out}}(r, \theta) = \frac{1}{4\pi\epsilon_0} \left(\sum_{l=0}^{\infty} \frac{P_l(\cos\theta)}{d^{l+1}} \left(\frac{q r^l}{r} + q' \frac{a^{2l} d}{r^{l+1}} \right) \right) \quad |\vec{x}| < d$$

Inside The effects are not quite that of a pt. charge q at d , but there is no reason to move the pt. charge. Place q'' at d instead.

$$\rightarrow \text{Then } \Phi_{\text{in}} = \frac{1}{4\pi\epsilon_0} \left(\frac{q''}{|\vec{x}-d\hat{y}|} \right) = \frac{q''}{4\pi\epsilon_0} \sum_{l=0}^{\infty} \frac{r^l}{d^{l+1}} P_l(\cos\theta) \quad \checkmark$$

Now all we have to do is find q' and q'' in terms of q

Boundary conditions to the rescue

- ① On surface: $(\vec{D}_2 - \vec{D}_1) \cdot \hat{n} = \sigma$
- ② Tangential components are equal: $\vec{E}_1 \cdot \hat{\theta} = \vec{E}_2 \cdot \hat{\theta}$

① $\vec{D} = \epsilon \vec{E}$, set 2 \rightarrow out, 1 \rightarrow in, and $\sigma = 0$ b/c no charge on surface.

$$(\epsilon_0 \vec{E}_{\text{out}} - \epsilon \vec{E}_{\text{in}}) \cdot \hat{r} = 0 \Rightarrow \epsilon_0 \vec{E}_{\text{out}} \cdot \hat{r} = \epsilon \vec{E}_{\text{in}} \cdot \hat{r} \Rightarrow \left. \frac{\partial \Phi_{\text{out}}}{\partial r} \right|_{r=a} = \left. \frac{\partial \Phi_{\text{in}}}{\partial r} \right|_{r=a}$$

We can do this term by term, since summations are the same on both sides.

$$\left. \frac{\epsilon}{4\pi\epsilon_0} \frac{q''}{d^{l+1}} \frac{r^{l-1}}{r^{l+1}} P_l(\cos\theta) \right|_{r=a} = \left. \left(\frac{q}{4\pi\epsilon_0} \frac{r^{l-1}}{d^{l+1}} P_l(\cos\theta) - \frac{q'}{4\pi\epsilon_0} \frac{r^{l-2+2l}}{d^l} P_l(\cos\theta) \right) \right|_{r=a}$$

$$\frac{\epsilon}{\epsilon_0} \frac{q''}{d^{l+1}} \frac{a^{l-1}}{a^{l+1}} = \left(\frac{q}{\epsilon_0} \frac{a^{l-1}}{d^{l+1}} - \frac{q'}{\epsilon_0} \frac{d a^{2l}}{r^{l+2}} - \frac{q'}{\epsilon_0} \frac{d a^{2l}}{r^{l+2}} \right) \frac{\epsilon_0}{\epsilon_0} \Big|_{r=a}$$

$$q'' a^{l-1} = q a^{l-1} - q' \frac{d a^{2l}}{r^{l+2}} - q' \frac{d a^{2l}}{r^{l+2}} \Big|_{r=a} \Rightarrow q'' a^{l-1} = q a^{l-1} - q' \frac{d}{a} a^{l-2} - q' d a^{l-2}$$

$$\Rightarrow q'' = q - q' \frac{d}{a} - q' \frac{d}{a} = q - q' \left(\frac{d}{a} + \frac{d}{a} \right)$$

② $\left. \frac{\partial \Phi_{\text{in}}}{\partial \theta} \right|_{r=a} = \left. \frac{\partial \Phi_{\text{out}}}{\partial \theta} \right|_{r=a} \Rightarrow$ (term by term) $\frac{q''}{\epsilon} \frac{a^l}{d^{l+1}} = \frac{1}{\epsilon_0 d^{l+1}} (q a^l + q' a^{l-1} d)$

$$\frac{q''}{\epsilon} = \frac{1}{\epsilon_0} \left(q + \frac{q' d}{a} \right)$$

Solve these two equations simultaneously for q' and q'' (see Mathematica).

$$\left\{ q' = \frac{-a l q (\epsilon - \epsilon_0)}{d (l \epsilon + \epsilon_0 + l \epsilon_0)} \right\}, \left\{ q'' = \frac{(1 + 2l) q \epsilon}{\epsilon \epsilon_0 + \epsilon_0 + l \epsilon_0} \right\} \quad \checkmark$$

(b) Near the center of the sphere, r is small, so $r \ll d$, and $\frac{r}{d} \ll 1$

so write $\Phi_{in} = \frac{q''}{4\pi\epsilon d} \sum_{\ell=0}^{\infty} \left(\frac{r}{d}\right)^{\ell} P_{\ell}(\cos\theta) \approx \frac{q''}{4\pi\epsilon d} \left[1 + \frac{r}{d}(\cos\theta)\right]$ expand

$= \frac{(1+2\ell)q}{2\epsilon + \epsilon_0 + \ell\epsilon_0} \frac{1}{4\pi d} \left[1 + \frac{r}{d}(\cos\theta)\right]$ and note $r\cos\theta = y$ (z in standard spherical coordinates, but you're at liberty to do as you please!)

$= \frac{(1+2\ell)q}{2\epsilon + \epsilon_0 + \ell\epsilon_0} \frac{1}{4\pi d} \left[1 + \frac{y}{d}\right]$

and $\vec{E}_{in} = -\frac{\partial}{\partial y} \Phi_{in} \hat{y} = -\frac{(1+2\ell)q}{4\pi d^2 (2\epsilon + \epsilon_0 + \ell\epsilon_0)} \hat{y} = \frac{-q}{4\pi\epsilon_0 d^2} \left[\frac{3}{\frac{\epsilon}{\epsilon_0} + 2}\right] \hat{y}$ ✓

(c) Now deal with the limit $\frac{\epsilon}{\epsilon_0} \rightarrow \infty$

$\Phi_{in} = \frac{q}{4\pi\epsilon_0 d} \underbrace{\left(\frac{1+2\ell}{2\frac{\epsilon}{\epsilon_0} + 1 + 2\ell}\right)}_{\rightarrow 1} \sum_{\ell=0}^{\infty} \left(\frac{r}{d}\right)^{\ell} P_{\ell}(\cos\theta)$

all terms go to zero except $\ell=0$. Nice!

$\Phi_{in, \epsilon/\epsilon_0 \rightarrow \infty} = \frac{q}{4\pi\epsilon_0 d} \left(\frac{1}{1}\right) \left(\frac{r}{d}\right)^0 = \frac{q}{4\pi\epsilon_0 d}$ ✓

and since $\vec{E}_{in, \epsilon/\epsilon_0 \rightarrow \infty} = -\vec{\nabla} \Phi_{in, \epsilon/\epsilon_0 \rightarrow \infty} = \vec{0}$, we get the same result as that for a conducting sphere. It shields the inside from the electric field, and maintains constant potential due completely to external charge.

for Φ_{out} , $\frac{1}{\epsilon/\epsilon_0 \rightarrow \infty} = \frac{1}{d\ell} q\ell a = \frac{q a}{d}$

Thus $\Phi_{out, \epsilon/\epsilon_0 \rightarrow \infty}(r, \theta) = \frac{q}{4\pi\epsilon_0} \left(\sum_{\ell=0}^{\infty} \frac{P_{\ell}(\cos\theta)}{r^{\ell+1}} \left(d^{\ell} + \frac{a^{2\ell+1}}{d^{\ell+1}} \right) \right)$

$= \frac{q}{4\pi\epsilon_0} \left(\frac{a/d}{r} + \frac{1}{|r-d|} - \frac{a/d}{|r-a^2/d|} \right)$ ✓

Just as we have done 1 Million times with normal image charges.

(10)

Mathematica for part a

I use $q'' = q_2$ and $q' = q_1$

In[109]:= Solve[{{q2 == q - q1 (d/(L*a) + d/a), q2/epsilon == 1/epsilon0 (q + q1*d/a)}, {q1, q2}]

Out[109]:= {{q1 -> -a L q (epsilon - epsilon0) / (d (L epsilon + epsilon0 + L epsilon0)), q2 -> (1 + 2 L) q epsilon / (L epsilon + epsilon0 + L epsilon0)}}