

4 Jackson 3.20

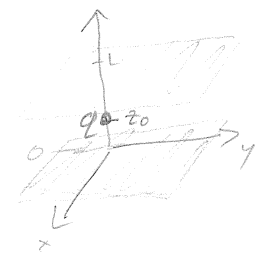
a) $G(\vec{x}, \vec{x}') = \frac{1}{L} \sum_{n=1}^{\infty} \sum_{m=-\infty}^{\infty} e^{im(\phi-\phi')} \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right) I_m\left(\frac{n\pi}{L} \rho_c\right) K_m\left(\frac{n\pi}{L} \rho_s\right)$

The potential can again be found...

$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int \rho(\vec{x}') G(\vec{x}, \vec{x}') d^3x' - \frac{1}{4\pi} \oint_S \Phi(\vec{x}') \frac{\partial G}{\partial n'} da'$

The charge Distribution is

$\rho(\vec{x}') = \rho_s \delta(z'-z_0) \delta(\rho')$ for a pt. charge @ z_0 . ($\Phi(\vec{x}') = 0$ is given)



(invariant in ϕ)
Integrate to get A:

$\int \rho(\vec{x}') dV' = q = A \int_0^{2\pi} \int_0^{\infty} \int_{-\infty}^{\infty} \delta(z'-z_0) \delta(\rho') \rho' dz' d\rho' d\phi' = A 2\pi \rho_s \Rightarrow A = \frac{q}{2\pi \rho_s}$

Thus $\rho(\vec{x}') = \frac{q}{2\pi \rho_s} \delta(z'-z_0) \delta(\rho')$

Set the charge to be on the z-axis. Then $\rho' = 0$.

Obviously then ρ' is ρ_c (so ρ must be ρ_s).

$I_m\left(\frac{n\pi}{L} \rho_c\right) = I_m(0) = 0$ for all $m \neq 0$.

Now we can integrate $\rho(\vec{x}')$ w/ $G(\vec{x}, \vec{x}')$ to get $\Phi(\vec{x})$ which is what we seek.

$$\Phi(\vec{x}) = \Phi(\rho, z) = \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^{\infty} \int_{-\infty}^{\infty} d\phi' dz' \rho' \rho \left(\frac{q}{2\pi \rho_s} \delta(z'-z_0) \delta(\rho') \right) \frac{1}{L} \sum_{n=1}^{\infty} \sum_{m=-\infty}^{\infty} \left[e^{im(\phi-\phi')} \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right) I_m\left(\frac{n\pi}{L} \rho_c\right) K_m\left(\frac{n\pi}{L} \rho_s\right) \right]$$

$$= \frac{q}{2\pi^2 \epsilon_0 L} \sum_{n=1}^{\infty} \sum_{m=-\infty}^{\infty} \left[\sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z_0}{L}\right) K_m\left(\frac{n\pi}{L} \rho_s\right) \int_0^{2\pi} \int_0^{\infty} \delta(\rho') e^{im(\phi-\phi')} d\phi' d\rho' \right] K_m\left(\frac{n\pi}{L} \rho_s\right)$$

$$= \frac{q}{2\pi^2 \epsilon_0 L} \sum_{n=1}^{\infty} \sum_{m=-\infty}^{\infty} \left[\sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z_0}{L}\right) K_m\left(\frac{n\pi}{L} \rho_s\right) e^{im\phi} \int_0^{2\pi} e^{-im\phi'} d\phi' \right]$$

Now, if $m=0$, $\int_0^{2\pi} e^{-im\phi'} d\phi' = \int_0^{2\pi} d\phi' = 2\pi$
 if $m \neq 0$, $\int_0^{2\pi} e^{-im\phi'} d\phi' = \frac{1}{-im} e^{-im\phi'} \Big|_0^{2\pi} = \frac{1}{-im} e^{-im2\pi} - \frac{1}{-im} = \frac{1}{-im} - \frac{1}{-im} = 0$
 Thus $m=0$ is the only case that matters.

$$\Phi(\rho, z) = \frac{q}{\pi \epsilon_0 L} \sum_{n=1}^{\infty} \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z_0}{L}\right) K_0\left(\frac{n\pi}{L} \rho_s\right)$$
 Azimuthally symmetric, as expected.

b) $\sigma_o(\rho) = -\epsilon_0 \frac{\partial \Phi}{\partial z} \Big|_{z=0} = \frac{q}{\pi L} \sum_{n=1}^{\infty} \frac{\partial}{\partial z} \left(\sin\left(\frac{n\pi z}{L}\right) \right) \Big|_{z=0} \sin\left(\frac{n\pi z_0}{L}\right) K_0\left(\frac{n\pi}{L} \rho_s\right)$

$$\rightarrow \frac{n\pi}{L} \cos\left(\frac{n\pi z}{L}\right) \Big|_{z=0} = \frac{n\pi}{L} \cos(0) = \frac{n\pi}{L}$$

$$\sigma_o(\rho) = \frac{q}{L^2} \sum_{n=1}^{\infty} n \sin\left(\frac{n\pi z_0}{L}\right) K_0\left(\frac{n\pi}{L} \rho_s\right)$$

$\sigma_L(\rho) = -\epsilon_0 \frac{\partial \Phi}{\partial z} \Big|_{z=L} = \frac{q}{\pi L} \sum_{n=1}^{\infty} \frac{\partial}{\partial z} \left(\sin\left(\frac{n\pi z}{L}\right) \right) \Big|_{z=L} \sin\left(\frac{n\pi z_0}{L}\right) K_0\left(\frac{n\pi}{L} \rho_s\right)$

$$\sigma_L(\rho) = \frac{q}{L^2} \sum_{n=1}^{\infty} n (-1)^n \sin\left(\frac{n\pi z_0}{L}\right) K_0\left(\frac{n\pi}{L} \rho_s\right) \rightarrow \frac{n\pi}{L} \cos\left(\frac{n\pi z}{L}\right) \Big|_{z=L} = \frac{n\pi}{L} \cos(n\pi) = \frac{n\pi}{L} (-1)^n$$

Comparison to results from 3.19?