

40.3

To derive the given expression for $\frac{dP_m}{d\Omega}$, we begin w/ eq 38.42:

$$\frac{dP_m}{d\Omega} = \frac{\omega_0 e^2}{2\pi R} \beta^3 m^2 \left([J'_m(m\beta \sin\theta)]^2 + \left[\frac{J_m(m\beta \sin\theta)}{\beta \tan\theta} \right]^2 \right)$$

By 38.53, $J_m(z) = \int_0^\pi \frac{1}{\pi} d\phi \cos(z \sin\phi - m\phi)$

$\Rightarrow J_m(m\beta \sin\theta) = \int_0^\pi \frac{1}{\pi} d\phi \cos(m\beta \sin\theta \sin\phi - m\phi)$

\checkmark We assume $\omega_0 r = \phi$ is small, so we can Taylor expand $\sin\phi \approx \phi - \frac{\phi^3}{6}$

$\checkmark \Rightarrow J_m(m\beta \sin\theta) = \int_0^\pi \frac{1}{\pi} d\phi \cos(m\beta \phi \sin\theta - \frac{m\beta \sin\theta}{6} \phi^3 - m\phi)$

$$= \int_0^\pi \frac{1}{\pi} d\phi \cos(-m\phi(1 - \beta \sin\theta) - \frac{1}{6} m\beta \sin\theta \phi^3)$$

\checkmark • note $1 - \beta \sin\theta = \frac{1}{2}(1 - \beta \sin^2\theta) = \frac{1}{2}\epsilon$... (given in the problem)

\checkmark • note $\cos(-x) = \cos(x)$, so we flip the sign in the argument

$$= \int_0^\pi \frac{1}{\pi} d\phi \cos\left(\frac{m}{2}\phi\epsilon - \frac{1}{6} m\beta \sin\theta \phi^3\right)$$

• for the ultrarelativistic limit, $\theta \rightarrow \frac{\pi}{2}$, and $\beta \rightarrow 1$

Solution continues on the next page....

thus $\int J_m(m\beta \sin \theta) \approx \int_0^\pi \frac{1}{\pi} d\phi \cos\left(\frac{m}{2}(\phi \epsilon + \frac{\phi^3}{3})\right)$

- change of variables: $\psi = \phi \left(\frac{m}{2}\right)^{1/3}$

- in the ultrarelativistic limit The lower limit of integration stays at $\psi = 0 \left(\frac{m}{2}\right)^{1/3} = 0$

But, The upper limit goes to $\psi = (\pi) \left(\frac{m}{2}\right)^{1/3}$, which in the ultrarelativistic limit $\rightarrow \infty$.

$$J_m(m\beta \sin \theta) \approx \left(\frac{2}{m}\right)^{1/3} \int_0^\infty \frac{1}{\pi} d\psi \cos\left(\left(\frac{m}{2}\right)^{2/3} \epsilon \psi + \frac{1}{3} \psi^3\right)$$

- make the two following definitions:

$$\left(\frac{m}{2}\right)^{2/3} \epsilon \equiv z$$

$$t \equiv \psi$$

$$\int J_m(m\beta \sin \theta) \approx \frac{1}{\sqrt{3}} \sqrt{\frac{3\epsilon}{z}} \int_0^\infty \frac{1}{\pi} dt \cos\left(zt + \frac{t^3}{3}\right) = \frac{\epsilon}{\pi\sqrt{3}} K_{1/3}(y)$$

by the definition of $K_{1/3}$ in the problem.

- next, $z = \left(\frac{3y}{2}\right)^{2/3}$ from the problem statement, so

$$y = \frac{2z^{3/2}}{3} = \frac{2}{3} \left(\frac{m}{2}\right) \epsilon^{3/2} = \frac{m}{3} \epsilon^{3/2}$$

$$\Rightarrow J_m(m\beta \sin \theta) = \frac{\epsilon}{\pi\sqrt{3}} K_{1/3}\left(\frac{m}{3} \epsilon^{3/2}\right)$$

We plug this into The second term in parentheses of the original equation:

$$\left[\frac{J_m(m\beta \sin \theta)}{\beta \tan \theta} \right]^2 = \frac{\epsilon K_{1/3}^2 \left(\frac{m}{3} \epsilon^{3/2}\right)}{\pi^2 \beta^2 \tan^2 \theta} \approx \frac{\epsilon \cos^2 \theta K_{1/3}^2 \left(\frac{m}{3} \epsilon^{3/2}\right)}{\pi^2 \beta}$$

where I sent $\beta \rightarrow 1$, and $\tan \theta \approx \sin \theta = \frac{1}{\cos \theta}$ in the small angle approx.

Other term (very similar, so I'll go faster):

$$J'_m(z) = \frac{d}{dz} \int_0^\pi \frac{1}{\pi} d\phi \cos(z \sin \phi - m\phi) = - \int_0^\pi \frac{d\phi}{\pi} \sin \phi \sin(z \sin \phi - m\phi)$$

small angle approx $\approx - \int_0^\pi \frac{d\phi}{\pi} \phi \sin(z \sin \phi - m\phi)$

Using the same argument as before,

$$J'_m(m\beta \sin \theta) \approx \int_0^\pi \frac{d\phi}{\pi} \phi \sin\left(\frac{m}{2}(\epsilon \phi + \frac{\phi^3}{3})\right)$$

again, let $\psi = \phi \left(\frac{m}{2}\right)^{1/3}$, and the upper limit of integration $\left(\frac{m}{2}\right)^{1/3} \pi \rightarrow \infty$

$$\Rightarrow J'_m(m\beta \sin \theta) = \left(\frac{2}{m}\right)^{2/3} \int_0^\infty \frac{1}{\pi} d\psi \sin\left[\left(\frac{m}{2}\right)^{2/3} \epsilon \psi + \frac{\psi^3}{3}\right]$$

now $\psi \rightarrow t$ and $\left(\frac{m}{2}\right)^{2/3} \epsilon \rightarrow z$

$$= \frac{1}{\pi} \frac{\epsilon}{z} \int_0^\infty dt t \sin\left(tz + \frac{t^3}{3}\right) = \frac{\epsilon}{\pi\sqrt{3}} \frac{\sqrt{3}}{z} \int_0^\infty dt t \sin\left(tz + \frac{t^3}{3}\right)$$

$$= \frac{\epsilon}{\pi\sqrt{3}} K_{2/3}(y) = \frac{\epsilon}{3\pi} K_{2/3}\left(\frac{m}{3} \epsilon^{3/2}\right)$$

$$\Rightarrow \left[J'_m(m\beta \sin \theta) \right]^2 \approx \frac{\epsilon^2}{3\pi^2} K_{2/3}^2\left(\frac{m}{3} \epsilon^{3/2}\right)$$

$$\Rightarrow \frac{dP_m}{d\Omega} = \frac{\omega_0 e^2 m^2}{2\pi R} \beta^{2\gamma-1} \left[\frac{\epsilon^2}{3\pi^2} K_{2/3}^2\left(\frac{m}{3} \epsilon^{3/2}\right) + \frac{\epsilon}{\pi^2 \beta} \cos^2 \theta K_{1/3}^2\left(\frac{m}{3} \epsilon^{3/2}\right) \right]$$

$$\Rightarrow \left[\frac{dP_m}{d\Omega} = \omega_0 \frac{e^2 m^2}{R 6 \pi^3} \left[\epsilon^2 K_{2/3}^2\left(\frac{m}{3} \epsilon^{3/2}\right) + \epsilon \cos^2 \theta K_{1/3}^2\left(\frac{m}{3} \epsilon^{3/2}\right) \right] \right]$$