

Sakurai 1.29

@ $[x_i, G(\vec{p})] = [x_i, G(p_i)]$ - because we can always consolidate the other terms (p_j, p_k, \dots) will always commute w/ x_i , and those terms will go to zero. \rightarrow They'll stick around, but they will not be affected by differentiation with respect to x_i .

Taylor series:

$$G(p_i) = \sum_{n=0}^{\infty} g_n p_i^n$$

Note:

$$[A, B^n] = n B^{n-1} [A, B] \rightarrow \text{where did this come from? (can be proved for } x_i \text{ and } p_i \text{ by induction)}$$

$$\begin{aligned} [x_i, G(\vec{p})] &= [x_i, \sum_{n=0}^{\infty} g_n p_i^n] = [x_i, g_1 p_i] + [x_i, g_2 p_i^2] + \sum_{n=3}^{\infty} [x_i, g_n p_i^n] \\ &= \sum_{n=0}^{\infty} g_n [x_i, p_i^n] = \sum_{n=0}^{\infty} g_n n p_i^{(n-1)} [x_i, p_i] \quad \leftarrow \text{see above identity.} \\ &= i\hbar \sum_{n=0}^{\infty} n g_n p_i^{(n-1)} = i\hbar \frac{\partial G}{\partial p_i} \quad \checkmark \quad \square \end{aligned}$$

(because $\frac{\partial}{\partial p_i} G = \frac{\partial}{\partial p_i} \sum_{n=0}^{\infty} g_n p_i^n = \sum_{n=0}^{\infty} g_n n p_i^{n-1}$)

Likewise

$$\begin{aligned} [p_i, F(\vec{x})] &= [p_i, F(x_i)] = [p_i, \sum_{n=0}^{\infty} f_n x_i^n] = \sum_{n=0}^{\infty} f_n [p_i, x_i^n] = \sum_{n=0}^{\infty} f_n n x_i^{n-1} [p_i, x_i] \\ &= -i\hbar \sum_{n=0}^{\infty} f_n n x_i^{n-1} = -i\hbar \frac{\partial F}{\partial x_i} \quad \checkmark \quad \square \end{aligned}$$

as before

$$\begin{aligned} \textcircled{b} [x^2, p^2] &= [x^2, p]p + p[x^2, p] = x[x, p]p + [x, p]xp \\ &\quad + px[x, p] + p[x, p]x \\ &= x i\hbar p + i\hbar xp + px i\hbar + p i\hbar x \\ &= i\hbar(2xp + 2px) = \boxed{2i\hbar(xp + px)} \quad \checkmark \end{aligned}$$

$$[x^2, p^2]_{\text{classical}} = \frac{\partial x^2}{\partial x} \frac{\partial p^2}{\partial p} - \frac{\partial x^2}{\partial p} \frac{\partial p^2}{\partial x} = 2x \cdot 2p - 0 = 4xp \quad \checkmark$$

This is by 1.6.48 which states:

$$[A(q, p), B(q, p)]_{\text{classical}} = \sum_s \left(\frac{\partial A}{\partial q_s} \frac{\partial B}{\partial p_s} - \frac{\partial A}{\partial p_s} \frac{\partial B}{\partial q_s} \right)$$

(Poisson brackets! You'll encounter these in classical mechanics soon.)
 where $A \rightarrow x^2$, $B \rightarrow p^2$, $q \rightarrow x$, $p \rightarrow p$ (duh),
 and we have just one s .