Pathia 2-7 (1) The first goal is to determine an "asymptotic" expression for J2 N I-D SHO's, with energies $E(n_r) = (n_r + \frac{1}{2})\hbar \omega_0$
Total energy of system $E = \frac{E \epsilon(n_r)}{E} = \frac{\sum (n_r + \frac{1}{k}) \overline{h} \omega_o}{\sum (n_r + \frac{1}{k}) \overline{h} \omega_o} = \frac{\overline{h} \omega_o \overline{\sum n_r} + \frac{11}{k} \overline{h} \omega_o}{\sum (n_r + \frac{1}{k}) \overline{h} \omega_o}$
$\frac{2n_1}{}$ = $\frac{2}{n\omega_0}$ + $\frac{2}{3}$ = R. Wen $E \gg N$ ($R \propto \frac{E}{\hbar \omega_0}$) 12 is the number of ways to sum up the n_r 's so that thy equal R .
ing kara Perengan sa Tinggalang Kalimatan (Pangangan Sangangan) na Ka

Scroll Down **—** *Solution continued on next page*

Following the procedure I used in question 1-8, we use 3.825;
 $DE = (R+N-1)!$ \Rightarrow $ln D = Rln(R+1) + Nln(R+1)$
 $R!(N-1)!$ $2 - \lambda$, com. as before, where R, N >> I. Furthon, in this book "assumptivally" meas $R >> N$, so we have $\frac{R}{N}$
 $\frac{ln 52 \approx Rln(R) + Nln(R) = ln(R^0) \approx ln(R^0) \times ln(R^1) \times ln(R^1)}{o(1 - N!)}$ Finally, $D = R^N$ Now plug in the circled Same $\left(\frac{E/\hbar\omega_0}{N}\right)^N$
N! expression for R on provious page => (ii) We begin by witting the Hamiltonian of a single hamonic oscillato:
 $E = P^2 + q^2$
 $2m \frac{2m\omega_e^2}{2m\omega_e^2}$ Thus, for N oscillators we have $E = \sum_{r=1}^{n} \frac{P_r^2}{2m} + \frac{q_r^2}{2m}$ As you may have vioticed, this (1) is just the equation of a of this phase space enclosed with total energy less than E.
Though it would be very nice, I don't know how
to integrate over this volume, so lets transform to a hypersphere.
 $let P'_i = Pr \quad end P'_i = q_v \quad (so dp'_i = (\frac{1}{2m})^{w_i} dp_r \cdot dq_i = (\frac{1}{$ $\frac{10}{10}$ $\frac{N}{\sigma w}, E = \sum_{i=1}^{N} (P_i^2 + q_i^2) \cup \rightarrow \infty$ hyperspine of radius \sqrt{E} $V(E) = \int \cdots d^{N}q d^{N}p \longrightarrow \cdots \qquad d^{N}q \cdot d^{N}p \longrightarrow \frac{d^{N}q}{m\omega_{o}^{2}}$
 $O \leq \frac{1}{k+1} \leq E$
 $O \leq \frac{1}{k$ This allows us now to use φ , $\vec{\tau}$, appendix C .
 $V(E) = \left(\frac{4m}{m\omega_o^2}\right)^{N/2} \overline{1}^{\frac{2N/2}{2N/2}} \overline{E}^{\frac{2N}{2N}} = \left(\frac{2}{\omega_o}\right)^N \overline{1}^N \overline{E}^N$ $\frac{1}{\sqrt{1-\frac{1$ Now we note $V = C \Im 2$ where C is a constant. That constant is clearly h^N